

402(6) : Prece Definition of the Relativistic Lagrangian and Summary of Newtonian Mechanics

The Cartesian coordinates of Lagrangian is defined by :

$$L = -mc^2 \left(1 - \left(\frac{v_x^2 + v_y^2 + v_z^2}{c^2} \right)^{1/2} \right) - U \quad (1)$$

where

$$v^2 = v_x^2 + v_y^2 + v_z^2 \quad (2)$$

is the speed of one reference frame with respect to another. The relativistic momentum is defined by :

$$p_i = \gamma m v_i \quad (3)$$

where

$$i = x, y, z \quad (4)$$

It follows that :

$$p_x = \frac{\partial L}{\partial v_x} \quad (5)$$

and so on, Q.E.D.

Summary of Newtonian Dynamics
The Newtonian orbital equation is :

$$\underline{F} = m \underline{\ddot{r}} = - \frac{m b}{r^3} \underline{r} \quad (6)$$

This was developed in Cartesian coordinates in UFT377. In plane polar coordinates :

$$\underline{r} = r \underline{e}_r \quad (7) \quad (8)$$

$$\underline{\ddot{r}} = (\ddot{r} - r\dot{\phi}^2) \underline{e}_r + (r\ddot{\phi} + 2\dot{r}\dot{\phi}) \underline{e}_\phi$$

From eqs. (1) to (8) the following two equations are obtained :

$$\ddot{r} - r\dot{\phi}^2 = - \frac{mb}{r^2} \quad (9)$$

and

$$r\ddot{\phi} + 2\dot{r}\dot{\phi} = 0 \quad (10)$$

2) Eq. (9) is the Leibnitz equation of orbits, and eq. (10) is the conservation of angular momentum. Eq. (9) is obtained from the Euler Lagrange equation:

$$\frac{dL}{dr} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) \quad (11)$$

with the Lagrangian:

$$\begin{aligned} L &= \frac{1}{2} m \underline{v} \cdot \underline{v} - U \\ &= \frac{1}{2} m \underline{\dot{r}} \cdot \underline{\dot{r}} + \frac{mMG}{r} \quad (12) \\ &= \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2) + \frac{mMG}{r} \end{aligned}$$

Eq. (10) is obtained from

$$\frac{\partial L}{\partial \phi} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) \quad (13)$$

The generalized momentum from eq. (11) is the linear momentum:

$$p = \frac{\partial L}{\partial \dot{r}} \quad (14)$$

The generalized momentum from eq. (13) is the angular momentum:

$$L = \frac{\partial L}{\partial \dot{\phi}} \quad (15)$$

It follows from eq. (13) that:

$$\frac{dL}{dt} = 0 \quad (16)$$

i.e. angular momentum is conserved.

2) The angular momentum is:

$$L = m r^2 \dot{\phi} \quad - (17)$$

It follows that

$$\dot{\phi} = \frac{L}{m r^2} \quad - (18)$$

$$\ddot{\phi} = \frac{d\dot{\phi}}{dt} = \frac{d}{dt} \left(\frac{L}{m r^2} \right)$$

$$= \frac{d}{dr} \left(\frac{L}{m r^2} \right) \frac{dr}{dt}$$

$$= - \frac{2 \dot{r} L}{m r^3} \quad - (19)$$

It follows from eqs. (18) and (19) that:

$$r \ddot{\phi} + 2 \dot{r} \dot{\phi} = 0 \quad - (20)$$

which is eq. (10), O.C.D.

The Newtonian Hamiltonian is:

$$H = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2) - \frac{m M G}{r} \quad - (21)$$

and is conserved:

$$\frac{dH}{dt} = 0 \quad - (22)$$

So, H and L are constants of motion.

Using eq. (17):

$$\begin{aligned} \frac{dH}{dt} &= \dot{r} \frac{dH}{dr} = \frac{1}{2} m \dot{r} \frac{d}{dr} \left(\dot{r}^2 + \frac{L^2}{m r^2} \right) + \dot{r} \frac{m M G}{r^2} \\ &= 0 \quad - (23) \end{aligned}$$

Using $\frac{d}{dr}(\dot{r}^2) = \frac{d}{dt}(\dot{r}^2) \frac{1}{\dot{r}} = 2 \frac{\dot{r}}{r} \ddot{r} = 2 \ddot{r} - (24)$

The Leibnitz equation (9) follows:

$$\ddot{r} - \frac{L^2}{m^2 r^3} = -\frac{mG}{r^2} - (25)$$

Therefore:

$$\frac{dH}{dt} = 0$$

→

$$\ddot{r} - \frac{L^2}{m^2 r^3} = -\frac{mG}{r^2}$$

←

$$\frac{\partial \mathcal{L}}{\partial r} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{r}}$$

$$\frac{dL}{dt} = 0$$

→

$$r\ddot{\phi} + 2\dot{r}\dot{\phi} = 0$$

←

$$\frac{\partial \mathcal{L}}{\partial \phi} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\phi}}$$

The Leibnitz equation (9) may be developed as the Binet equation:

$$\frac{d^2}{d\phi^2} \left(\frac{1}{r} \right) + \frac{1}{r} = \frac{1}{d} - (26)$$

Let us take half right side rule is:

$$d = \frac{L^2}{m^2 mG} - (27)$$

By inspection, the conic section:

$$r = \frac{d}{1 + \epsilon \cos \phi} - (28)$$

solves the Binet equation. So the Newtonian dynamics solve the conic section.